

FTCS

For $\frac{\partial \phi}{\partial t} = -u \frac{\partial \phi}{\partial x}$, write as

$$\frac{\phi_i^{\tilde{t}+1} - \phi_i^{\tilde{t}}}{\Delta t} = -u \frac{(\phi_{i+1}^{\tilde{t}} - \phi_{i-1}^{\tilde{t}})}{2 \Delta x}$$

or ① $\phi_i^{\tilde{t}+1} = \phi_i^{\tilde{t}} - C [\phi_{i+1}^{\tilde{t}} - \phi_{i-1}^{\tilde{t}}] / 2$; $C = \frac{u \Delta t}{\Delta x} \equiv$ Courant number

② Assume solution $\phi(x, t) = \hat{\phi}(k, \omega) e^{i(kx + \omega t)}$

where $x = n \Delta x$, $t = \tilde{t} \Delta t$, $\omega =$ frequency, $k =$ wavenumber
 $n = 0, \pm 1, \pm 2$, $\tilde{t} = 0, 1, 2$

so that the continuous (x, t) space is replaced by a grid of discrete points.

Denote:

$$\phi_i^{\tilde{t}+1} = \phi_i^{\tilde{t}} \phi_0^1; \quad \phi_{i+1}^{\tilde{t}} = \phi_i^{\tilde{t}} \phi_1^0; \quad \phi_{i-1}^{\tilde{t}} = \phi_i^{\tilde{t}} \phi_{-1}^0$$

Substitute ② into ① and cancel common $\phi_i^{\tilde{t}}$ factor, $\hat{\phi}$ factor

$$e^{i \omega \Delta t} = 1 - \frac{C}{2} [e^{i k \Delta x} - e^{-i k \Delta x}]$$

Using Euler's formula $e^{i\theta} = \cos\theta + i\sin\theta$, $e^{-i\theta} = \cos\theta - i\sin\theta$

$$e^{i \omega \Delta t} = 1 - \frac{C}{2} [2i \sin(k \Delta x)] = 1 - C [i \sin(k \Delta x)]$$

Due to the imprecise approximation for the derivatives, ω is decomposed into real and imaginary parts

$$\omega = \omega_r + i\omega_i$$

Therefore

$$e^{i\omega r t} = e^{i(\omega_r + i\omega_i) r t} = e^{i\omega_r r t} e^{-\omega_i r t} = e^{i\omega_r r t} \lambda$$

Where $\lambda = e^{-\omega_i r t} \equiv$ amplitude change of the finite difference solution per time step.

In the true solution, $\lambda = 1$, but in a finite difference scheme, often $\lambda \neq 1$.

To solve for λ , expand $e^{i\omega_r r t}$ using Euler's formula

$$\lambda [\cos(\omega_r r t) + i \sin(\omega_r r t)] = 1 - C [i \sin(k a x)]$$

The real and imaginary parts must separately be equal:

$$\textcircled{3} \quad \lambda \cos(\omega_r r t) = 1$$

$$\textcircled{4} \quad \lambda \sin(\omega_r r t) = -C \sin(k a x)$$

$\textcircled{3}^2 + \textcircled{4}^2$ gives

$$\lambda^2 = 1 + C^2 \sin^2(k a x)$$

$$\lambda = \pm \sqrt{1 + C^2 \sin^2(k a x)}$$

Except for ~~small~~ very small k ($L \rightarrow \infty$), or for $\Delta x \rightarrow 0$, or for a 2 π wave, $|\lambda| > 1$.

FTCS is linearly unstable!

FTFS (for $U_i > 0$) , or forward in time, upstream

Will have

$$\Phi_i^{n+1} = \Phi_i^n - Q (\Phi_i^n - \Phi_{i-1}^n) ;$$

Linear analysis gives

$$e^{i\omega n \Delta t} = 1 - Q (1 - e^{-ik\Delta x})$$

or, applying λ and Euler's formula

$$\lambda (\cos(\omega_r \Delta t) + i \sin(\omega_r \Delta t)) = 1 - Q + Q \cos(k\Delta x) - Q i \sin(k\Delta x)$$

Equating real and imaginary parts

$$\textcircled{1} \quad \lambda \cos(\omega_r \Delta t) = 1 - Q + Q \cos(k\Delta x) = 1 - Q (1 - \cos(k\Delta x))$$

$$\textcircled{2} \quad \lambda \sin(\omega_r \Delta t) = -Q \sin(k\Delta x)$$

$$\text{or} \quad \lambda^2 = [1 - Q + Q \cos^2(k\Delta x)]^2 + [Q \sin(k\Delta x)]^2$$

$$\lambda^2 = 1 + 2Q(1 - \cos(k\Delta x)) + Q^2(1 - 2\cos(k\Delta x) + \cos^2(k\Delta x)) + Q^2 \sin^2(k\Delta x)$$

$$\lambda^2 = 1 + 2Q - 2Q \cos(k\Delta x) + Q^2 - 2Q^2 \cos(k\Delta x) + Q^2$$

$$\lambda^2 = 1 + 2Q [1 - \cos(k\Delta x) + Q^2 - 2Q \cos(k\Delta x)]$$

Turns out this is a damping scheme, with "mild" damping for small Q ($Q < 0.01$), strong damping for middle values of Q (0.1 to 0.4) with very high at 0.5)

preserved amplitude for $q=1.0$, and amplification for $q>1$. The damping is true for all $n \neq 2$, but is greatest at $n=2$.

Phase speed

To get analytic phase speed c , ⁽²⁾/₍₁₎

$$\tan(\omega_{rst}) = \frac{(2)}{(1)}$$

$$\text{Since } \omega = -ck = -c \frac{2\pi}{L_x} = -c \frac{2\pi}{n\lambda_x}$$

$$-c \frac{2\pi}{n\lambda_x} \Delta t = \tan^{-1}\left(\frac{(2)}{(1)}\right)$$

The true phase speed is $c_{true} = U$, so to compare the analytic c and true c , divide by U

$$\frac{c}{U} = \frac{-n}{2\pi} \frac{\Delta x}{U \Delta t} \tan^{-1}\left(\frac{(2)}{(1)}\right)$$

$\frac{1}{q}$

$$\frac{c_{anal}}{U} = -\frac{n}{2\pi q} \tan^{-1}\left(\frac{(2)}{(1)}\right)$$

$|c_{anal}|/U = 1.0$ at $q=0.5, 1$. For $q < 0.5$, c_{anal} too slow, and it is worse for small n ($=0$ for $n=2!$). For $q > 0.5$, c_{anal} is too fast, and it is worse for small n .